

Now  $\underline{f}^{(2)} = \underline{r}(s) + \underline{I}_\alpha \underline{\tilde{d}}_\alpha(s)$  as on p. 123,

But now  $\underline{\tilde{d}}_\alpha = \underline{R} \underline{\tilde{e}}_\alpha \equiv \underline{R} \underline{Q} \underline{e}_\alpha$ .

The same calculations as on pp. 133-134 yield

$$\begin{aligned} \underline{f}^{(2)} &= \underline{\tilde{d}}_\alpha \otimes \underline{e}_\alpha + (\underline{r}' + \underline{I}_\alpha \underline{\tilde{d}}_\alpha) \otimes \underline{e}_3 \\ &= \underline{R} \underline{Q} \underline{e}_\alpha \otimes \underline{e}_\alpha + (\underline{r}' + \underline{I}_\alpha \underline{R}' \underline{Q} \underline{e}_\alpha) \otimes \underline{e}_3 \end{aligned}$$

Summary  $\underline{R}, \underline{R}' \xrightarrow{\text{(1)}} \underline{R} \underline{Q}, \underline{R}' \underline{Q} \xrightarrow{\text{(2)}}$

Observe that (\*) p. 151 is equivalent to

$$\begin{aligned} \int_{\Omega} W(\underline{R} \underline{e}_\alpha \otimes \underline{e}_\alpha + (\underline{r}' + \underline{I}_\alpha \underline{R}' \underline{e}_\alpha) \otimes \underline{e}_3) d\underline{X}_1 d\underline{X}_2 \\ = \int_{\Omega} W(\underline{R} \underline{Q} \underline{e}_\alpha \otimes \underline{e}_\alpha + (\underline{r}' + \underline{I}_\alpha \underline{R}' \underline{Q} \underline{e}_\alpha) \otimes \underline{e}_3) d\underline{X}_1 d\underline{X}_2. \end{aligned}$$

where  $\underline{\hat{X}} = \underline{I}_\alpha \underline{e}_\alpha$  on the left and  $\underline{\hat{X}} = \underline{I}_\alpha \underline{Q} \underline{e}_\alpha$  on the right.  $\Rightarrow \underline{Q}(\Omega) = \Omega$  is a necessary condition for equality above — otherwise the integral on the right may not even make sense!

Next we want to realize (\*) p. 149 in the reduced form (\*) p. 148: Since we have "opened up the door" to  $\underline{R} \rightarrow \underline{R} \underline{Q}$  with  $\det(\underline{R} \underline{Q}) = \det \underline{R} \det \underline{Q}^{-1} = -1$ , we need to be

careful in extracting an axial vector as at the bottom of p. 141:

Lemma The axial vector of  $\underline{Q} \underline{W} \underline{Q}^T$  for  $\underline{Q} \in O(3)$ ,  $\underline{W} \in \text{Skew}(\mathbb{F}^3)$ , is given by

$$(\det \underline{Q}) \underline{Q} \underline{w},$$

where  $\text{axial}(\underline{W}) = \underline{w}$ . (Note  $\det \underline{Q} = \pm 1$ )

Pf. From the Lemma on p. 27, we have

$$\underline{Q} \underline{W} \underline{a} = \underline{Q} (\underline{w} \times \underline{a}) = \det \underline{Q} (\underline{Q} \underline{w}) \times (\underline{Q} \underline{a}).$$

Choose  $\underline{a} = \underline{Q}^T \underline{b} \Rightarrow$

$$\underline{Q} \underline{W} \underline{Q}^T \underline{b} = (\det \underline{Q}) (\underline{Q} \underline{w}) \times \underline{b} \quad \forall \underline{b}. \quad \square$$

$\therefore$  (\*) p. 151 leads

$$\underline{\Psi}(\underline{r}', \underline{R}, \underline{R}') = \underline{\Psi}(\underline{R}^T \underline{r}', (\det \underline{R}) \underline{R}^T \underline{k})$$

Thus (\*) p. 152 is equivalent to

$$(*) \quad \underline{\Psi}(\underline{Q}^T \underline{R}^T \underline{r}', \underbrace{\det(\underline{R} \underline{Q})}_{\det \underline{R} \det \underline{Q}} \underline{Q}^T \underline{R}^T \underline{k}) = \underline{\Psi}(\underline{R}^T \underline{r}', \underline{R}^T \underline{k})$$

For real deformations of the rod,  $\det \underline{R} = 1$ .

Claim: If (\*) above holds for  $\underline{Q} \in O_c(2)$ , then it also holds for  $\underline{Q}^T$  (let  $\underline{R} \rightarrow \underline{R} \underline{Q}^T$ ).

Also, let  $\bar{v} \equiv (v_1, v_2, v_3)$ ,  $\bar{k} \equiv (k_1, k_2, k_3)$   
and  $\bar{Q} \equiv [Q_{ij}] \leftarrow \text{rel. } \{\underline{e}_1, \underline{e}_2, \underline{e}_3\}$ .

Then in view of p. 137 & the development above, (\*) p. 152 is equivalent to

$$(*) \quad \mathcal{I}(\bar{Q}\bar{v}, (\det \bar{Q})\bar{Q}\bar{k}) = \mathcal{I}(\bar{v}, \bar{k})$$

for  $\bar{Q} \in O_c(2)$  a transverse symmetry.

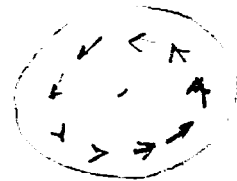
Two interesting cases:

$$\text{Symmetry group } SO_c(2) = \{ \underline{R}(\theta) : 0 \leq \theta < 2\pi \}$$

This is typically called hemitropy: (\*) above reads

$$\mathcal{I}(\underline{R}(\theta)\bar{v}, \underline{R}(\theta)\bar{k}) = \mathcal{I}(\bar{v}, \bar{k}) \quad \forall 0 \leq \theta < 2\pi$$

(\*\*) rotationally invariant  
but oriented



$$\begin{aligned} & \mathcal{I} \left( \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} (v_1, v_2), v_3, \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix} (k_1, k_2), k_3 \right) \\ &= \mathcal{I}(v_1, v_2, v_3, k_1, k_2, k_3) \quad \forall 0 \leq \theta < 2\pi \end{aligned}$$

Then (Cauchy - rotation invariance with more than one vector argument)

(\*\*)  $\Leftrightarrow$

$$\mathcal{I}(\underline{v}, \underline{k}) = \Gamma \left( v_1^2 + v_2^2, k_1^2 + k_2^2, v_1 k_1 + v_2 k_2, v_1 k_2 - v_2 k_1, \right. \\ \left. v_3, k_3 \right)$$

Observe: If  $\hat{\underline{v}} = v_1 \underline{e}_1 + v_2 \underline{e}_2$ ,  $\hat{\underline{k}} = k_1 \underline{e}_1 + k_2 \underline{e}_2$ ,  
then hemitropic

$$(*) \quad \mathcal{I}(\underline{v}, \underline{k}) = \Gamma \left( |\hat{\underline{v}}|^2, |\hat{\underline{k}}|^2, \hat{\underline{v}} \cdot \hat{\underline{k}}, \underline{e}_3 \cdot (\hat{\underline{v}} \times \hat{\underline{k}}), v_3, k_3 \right)$$

To quadratic order: assume straight state is "stress" free - 1<sup>st</sup> order terms vanish

$$\mathcal{I}(\underline{v}, \underline{k}) \doteq \frac{1}{2} \left\{ A |\hat{\underline{v}}|^2 + B |\hat{\underline{k}}|^2 + 2C \hat{\underline{v}} \cdot \hat{\underline{k}} \right.$$

$$(**) \quad \left. \begin{aligned} &+ 2D \underline{e}_3 \cdot (\hat{\underline{v}} \times \hat{\underline{k}}) + E (v_3 - 1)^2 + F k_3^2 \\ &+ 2G (v_3 - 1) k_3 \end{aligned} \right\}$$

$C, D, G \neq 0$  represent coupling coefficients,

In particular,  $G$  represents a coupling between twist  $k_3$  and axial strain  $v_3 - 1$ :

Proof (\*) p. 155  $\Rightarrow$  hemizopy (\*\*\*) p. 154  
is trivial.

$\leftarrow$

It suffices to show that

$$\mathcal{I}(\hat{\underline{v}}^*, \underline{v}_3, \hat{\underline{k}}^*, \underline{k}_3) = \mathcal{I}(\hat{\underline{v}}, \underline{v}_3, \hat{\underline{k}}, \underline{k}_3)$$

whenever,  $\hat{\underline{v}}^* \cdot \hat{\underline{v}}^* = \hat{\underline{v}} \cdot \hat{\underline{v}}, \hat{\underline{k}}^* \cdot \hat{\underline{k}}^* = \hat{\underline{k}} \cdot \hat{\underline{k}},$

$$\hat{\underline{v}}^* \cdot \hat{\underline{k}}^* = \hat{\underline{v}} \cdot \hat{\underline{k}}, \underline{e}_3 \cdot (\hat{\underline{v}}^* \times \hat{\underline{k}}^*) = \underline{e}_3 \cdot (\hat{\underline{v}} \times \hat{\underline{k}}).$$

Assume the latter conditions: The first two  $\Rightarrow$

$$\begin{aligned} \hat{\underline{v}}^* &= \underline{Q}_1 \hat{\underline{v}} & [Q_1] &= \begin{bmatrix} \hat{Q}_1 & \underline{0} \\ \underline{0} & 1 \end{bmatrix} \\ \hat{\underline{k}}^* &= \underline{Q}_2 \hat{\underline{k}} & [Q_2] &= \begin{bmatrix} \hat{Q}_2 & \underline{0} \\ \underline{0} & 1 \end{bmatrix}, \end{aligned}$$

$\hat{Q}_1, \hat{Q}_2$   $2 \times 2$  orthogonal matrices. Observe

$$\det \underline{Q}_\alpha = \det[\hat{Q}_\alpha] = \pm 1 \quad \alpha = 1, 2.$$

Then  $\hat{\underline{v}}^* \cdot \hat{\underline{k}}^* = \hat{\underline{v}} \cdot \underline{Q}_1^T \underline{Q}_2 \hat{\underline{k}} = \hat{\underline{v}} \cdot \hat{\underline{k}}$   
 $\Rightarrow \underline{Q}_1 = \underline{Q}_2 \equiv \underline{Q} \quad [Q] = \begin{bmatrix} \hat{Q} & \underline{0} \\ \underline{0} & 1 \end{bmatrix}.$

Finally  $\underline{e}_3 \cdot (\hat{\underline{v}}^* \times \hat{\underline{k}}^*) = \underline{e}_3 \cdot (\underline{Q} \hat{\underline{v}} \times \underline{Q} \hat{\underline{k}})$

( $\underline{Q} \underline{e}_3 = \underline{e}_3$ )  $= \underline{Q} \underline{e}_3 \cdot (\underline{Q} \hat{\underline{v}} \times \underline{Q} \hat{\underline{k}}) = \det \underline{Q} \underline{e}_3 \cdot (\hat{\underline{v}} \times \hat{\underline{k}}).$

$$\therefore \underline{e}_3 \cdot (\hat{\underline{v}}^* \times \hat{\underline{k}}^*) = \det \underline{Q} \underline{e}_3 \cdot (\hat{\underline{v}} \times \hat{\underline{k}})$$

$$\underline{e}_3 \cdot (\hat{\underline{v}} \times \hat{\underline{k}}) \Rightarrow \det \underline{Q} = \det [\hat{\underline{Q}}] = 1$$

proper  $\underline{Q} = \underline{R}(\Theta)$

$$\therefore \mathcal{I}(\hat{\underline{v}}^*, \nu_3, \hat{\underline{k}}^*, \kappa_3) = \mathcal{I}(\underline{R}(\Theta)\hat{\underline{v}}, \nu_3, \underline{R}(\Theta)\hat{\underline{k}}, \kappa_3)$$

$$= \mathcal{I}(\bar{\underline{R}}(\Theta)\bar{\underline{v}}, \bar{\underline{R}}(\Theta)\bar{\underline{k}})$$

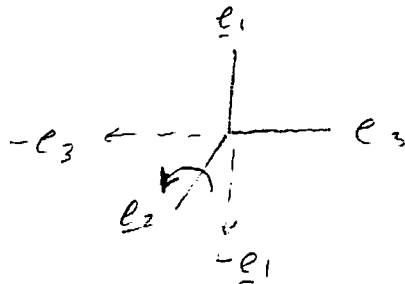
hemitropy

$$= \mathcal{I}(\bar{\underline{v}}, \bar{\underline{k}}) \quad \square$$

If the rod is homogeneous in the straight state (no explicit "s" dependence), then the rod enjoys another symmetry:

As on p. 152, in Experiment (3) we "flip" the rod over by performing a  $180^\circ$  rotation about, say,  $\underline{e}_2$ :  $\underline{G}\underline{e}_1 = -\underline{e}_1$ ,  $\underline{G}\underline{e}_2 = \underline{e}_2$

$$\underline{G}\underline{e}_3 = -\underline{e}_3$$



For simplicity, let's imagine doing this at the origin ( $s=0$ ), which need not be the true end of the rod.

The actual deformation we are performing is more complicated than described in the

course: the transformation is global:

$$\underline{f}(\underline{G}\underline{x}) = \underline{r}(-s) + \underline{x}_\alpha \underline{R}(-s) \underline{G} \underline{e}_\alpha,$$

ie.,  $\underline{r}(s) \rightarrow \underline{r}(-s)$

$$\underline{R}(s) = \underline{R}(-s) \underline{G}$$

$$\underline{G}^T = \underline{G}$$

$$\therefore \underline{R}_s^T \underline{r}'(s) \rightarrow -\underline{G}^T \underline{R}_s^T \underline{r}'(-s) = -\underline{G} \underline{R}^T \underline{r}'(-s)$$

$$\underline{R}^T \underline{K} \underline{R} \rightarrow \underline{G} \underline{R}^T (-\underline{R}' \underline{G} \underline{G} \underline{R}^T) \underline{R} \underline{G}$$

$$= -\underline{G} \underline{R}^T \underline{K} \underline{R} \underline{G} \quad (\det \underline{G} = 1)$$

$$\underline{R}^T \underline{K} \rightarrow -\underline{G} \underline{R}^T \underline{K}$$

(evaluate at  $s=0$ )

So the material symmetry reads:

$$\mathcal{F}(-\underline{G} \underline{R}^T \underline{v}, -\underline{G} \underline{R}^T \underline{k}) = \mathcal{F}(\underline{R}^T \underline{v}, \underline{R}^T \underline{k})$$

or  $\mathcal{F}(v_1, -v_2, v_3, k_1, -k_2, k_3) = \mathcal{F}(\underline{v}, \underline{k})$

Combining hemitropy + Mix symmetry (p.153):

$$v_1^2 + v_2^2 \rightarrow v_1^2 + (-v_2)^2 = v_1^2 + v_2^2$$

$$k_1^2 + k_2^2 \rightarrow k_1^2 + k_2^2$$

$$v_1 k_1 + v_2 k_2 \rightarrow v_1 k_1 + v_2 k_2$$