

Lecture 26

212

$$\frac{\partial}{\partial X_\alpha} \left(\frac{\partial \Phi}{\partial t_{i,\alpha}} \right) + \hat{b} = 0$$

$$\frac{\partial}{\partial X_\alpha} \left(\frac{\partial \Phi}{\partial d_{i,\alpha}} \right) - \frac{\partial \Phi}{\partial d} + \hat{l} = 0$$

These are static equil. equations - comparing with the static versions of (*) p. 198 and (*) p. 207, respectively, we deduce:

(*) $\hat{N}_\alpha = \frac{\partial \Phi}{\partial t_{i,\alpha}}, \quad \hat{\sigma} = \frac{\partial \Phi}{\partial d}, \quad \hat{M}_\alpha = \frac{\partial \Phi}{\partial d_{i,\alpha}}$

In anticipation of material-symmetry considerations, we can give (*) a more elegant tensorial form as follows:

A more accurate version of (*) p. 209 is

$$\Phi(\underline{t}, \underline{d}, \underline{d}_{i,\alpha}) = \Psi(\nabla_{\underline{r}}, \underline{d} \otimes \underline{e}_3, \nabla_{\underline{d}})$$

(recalling $\nabla_{\underline{r}} = t_{i,\alpha} \otimes \underline{e}_\alpha$, $\nabla_{\underline{d}} = d_{i,\alpha} \otimes \underline{e}_\alpha$).

Define $\underline{F} \equiv \nabla_{\underline{r}}, \quad \underline{G} \equiv \nabla_{\underline{d}}, \quad \underline{D} \equiv \underline{d} \otimes \underline{e}_3$.

Note $\underline{F}, \underline{G} \in L(\mathbb{E}^2, \mathbb{E}^3)$. Now consider

$$\frac{d}{d\varepsilon} \underline{\Phi}(\underline{r}_1 + \varepsilon \underline{n}, \underline{r}_2, \underline{d}, \underline{d}_1, \alpha) \Big|_{\varepsilon=0}$$

$$= \frac{d}{d\varepsilon} \underline{\Psi}((\underline{r}_1 + \varepsilon \underline{n}) \otimes \underline{e}_1 + \underline{r}_2, \underline{d} \otimes \underline{e}_3, \nabla \underline{d}) \Big|_{\varepsilon=0}$$

$$\frac{\partial \underline{\Phi}}{\partial \underline{r}_1} \cdot \underline{n} = \frac{\partial \underline{\Psi}}{\partial \underline{F}} \cdot \underline{n} \otimes \underline{e}_1 = \underline{n} \cdot \frac{\partial \underline{\Psi}}{\partial \underline{F}} \underline{e}_1$$

$$\Rightarrow \stackrel{(p.212)}{\underline{N}_1} = \frac{\partial \underline{\Phi}}{\partial \underline{r}_1} = \frac{\partial \underline{\Psi}}{\partial \underline{F}} \underline{e}_1 \quad \forall \underline{n} \in \mathbb{E}^3$$

p. 196

$$\Rightarrow \underline{N} = \frac{\partial \underline{\Psi}}{\partial \underline{F}}$$

Similarly $\underline{M} = \frac{\partial \underline{\Psi}}{\partial \underline{G}}$, and

$$\underline{\sigma} = \frac{\partial \underline{\Phi}}{\partial \underline{d}} = \frac{\partial \underline{\Psi}}{\partial \underline{D}} \underline{e}_3.$$

Material Objectivity

Borrowing again from the 3-d theory and (*) p. 209, we require the stored energy function to satisfy rotational invariance:

$$(M0) \quad \underline{\Phi}(\underline{Q} \underline{r}_\alpha, \underline{Q} \underline{d}, \underline{Q} \underline{d}_\alpha) = \underline{\Phi}(\underline{r}_\alpha, \underline{d}, \underline{d}_\alpha)$$

$$\forall \underline{Q} \in SO(3).$$

We now pursue a result analogously to that on p. 44:

Thm (M0) p. 213 + hyperelasticity p. 212
 \Rightarrow AMB p. 208 (presuming LMB and director momentum balance).
 * p. 207

PF As before, choose $\hat{Q}(t) = \exp(\underline{\Omega}t)$ with $\underline{\Omega} \in \text{skew}(\mathbb{E}^3)$. As shown on pp. 44-45, $\hat{Q}(t) \in \text{SO}(\mathbb{E}^3)$ for all t . Substituting $\hat{Q}(t)$ into (M0) and differentiation wrto t yields (using (*) p. 212):

$$\left[\underline{N}_\alpha \cdot \underline{\Omega} \exp(\underline{\Omega}t) \underline{r}_{i,\alpha} + \underline{\sigma} \cdot \underline{\Omega} \exp(\underline{\Omega}t) \underline{d} + \underline{M}_{i,\alpha} \cdot \underline{\Omega} \exp(\underline{\Omega}t) \underline{d}_{i,\alpha} \right] \Big|_{t=0} = 0 \quad \forall \underline{\Omega} \in \mathbb{E}^3$$

$$\Rightarrow \underline{N}_\alpha \cdot (\underline{\omega} \times \underline{r}_{i,\alpha}) + \underline{\sigma} \cdot (\underline{\omega} \times \underline{d}) + \underline{M}_{i,\alpha} \cdot (\underline{\omega} \times \underline{d}_{i,\alpha}) = 0$$

$$\forall \underline{\omega} = \text{axis}(\underline{\Omega}) \in \mathbb{E}^3$$

$$\begin{aligned} \text{Now } \underline{N}_\alpha \cdot (\underline{\omega} \times \underline{r}_{i,\alpha}) &= \underline{r}_{i,\alpha} \cdot (\underline{N}_\alpha \times \underline{\omega}) \\ &= \underline{\omega} \cdot (\underline{r}_{i,\alpha} \times \underline{N}_\alpha), \text{ etc.} \end{aligned}$$

$$\Rightarrow \underline{\omega} \cdot (\underline{r}_{i,\alpha} \times \underline{N}_\alpha + \underline{d} \times \underline{\sigma} + \underline{d}_{i,\alpha} \times \underline{M}_{i,\alpha}) = 0 \quad \forall \underline{\omega} \in \mathbb{E}^3$$

$$\Rightarrow \underline{r}_{i,\alpha} \times \underline{N}_\alpha + \underline{d} \times \underline{\sigma} + \underline{d}_{i,\alpha} \times \underline{M}_{i,\alpha} = 0. \quad \square$$

To exploit (M0) we write

$$\underline{\underline{F}}_m \equiv \nabla \underline{\underline{r}} + d\theta \underline{\underline{e}}_3 \equiv \underline{\underline{F}}_c \Big|_{\mathbb{R}_3=0}$$

mid-surface
deformation gradient,

and $\Psi(\underline{\underline{F}}_m, \underline{\underline{G}})$.

(M0) reads

$$(*) \quad \Psi(\underline{\underline{Q}} \underline{\underline{F}}_m, \underline{\underline{Q}} \underline{\underline{G}}) = \Psi(\underline{\underline{F}}_m, \underline{\underline{G}}) \quad \forall \underline{\underline{Q}} \in \text{SO}(\mathbb{E}).$$

Claim: (M0) (*) (above) \Rightarrow

$$\begin{aligned} \Psi(\underline{\underline{F}}_m, \underline{\underline{G}}) &= \mathcal{I}(\underline{\underline{F}}_m^T \underline{\underline{F}}_m, \underline{\underline{F}}_m^T \underline{\underline{G}}) \\ &\equiv \mathcal{I}(\underline{\underline{C}}_m, \underline{\underline{F}}_m^T \underline{\underline{G}}) \end{aligned}$$

Pf As before (p. 41) write $\underline{\underline{F}}_m = \underline{\underline{R}}_m \underline{\underline{U}}_m$
and choose $\underline{\underline{Q}} \equiv \underline{\underline{R}}_m^T$ in (*) above:

$$\begin{aligned} \Rightarrow \Psi(\underline{\underline{F}}_m, \underline{\underline{G}}) &= \Psi(\underline{\underline{U}}_m, \underline{\underline{R}}_m^T \underline{\underline{G}}) \\ &\equiv \tilde{\Psi}(\underline{\underline{U}}_m^2, \underline{\underline{U}}_m \underline{\underline{R}}_m^T \underline{\underline{G}}) \\ &\equiv \mathcal{I}(\underline{\underline{C}}_m, \underline{\underline{F}}_m^T \underline{\underline{G}}). \quad \square \end{aligned}$$

Strain Measures

$$\begin{aligned} \underline{F}_m^T \underline{F}_m &= (\underline{e}_3 \otimes \underline{d} + \underline{\nabla}_r^T) (\underline{\nabla}_r + \underline{d} \otimes \underline{e}_3) \\ &= \underline{\nabla}_r^T \underline{\nabla}_r + \underline{\nabla}_r^T \underline{d} \otimes \underline{e}_3 + \underline{e}_3 \otimes \underline{d} \underline{\nabla}_r \\ &\quad + \underline{e}_3 \otimes \underline{d} (\underline{d} \otimes \underline{e}_3) \end{aligned}$$

Observe $\underline{\nabla}_r^T \in L(\mathbb{E}^3, \mathbb{E}^2)$

$$\Rightarrow \underline{\nabla}_r^T \underline{\nabla}_r \in L(\mathbb{E}^3, \mathbb{E}^2)$$

Indeed, $\underline{\nabla}_r = \underline{r}_\alpha \otimes \underline{e}_\alpha$

$$\underline{\nabla}_r^T \underline{\nabla}_r = \underline{e}_\alpha \otimes \underline{r}_\alpha \underline{r}_\beta \otimes \underline{e}_\beta$$

$$= \underline{r}_\alpha \cdot \underline{r}_\beta \underline{e}_\alpha \otimes \underline{e}_\beta$$

$$= \underset{\nearrow}{a_{\alpha\beta}} \underline{e}_\alpha \otimes \underline{e}_\beta \equiv \underline{\underline{C}} \text{ "in-plane" right Cauchy-Green}$$

components of "1st fund. form"
in differential geometry.

$$\underline{\nabla}_r^T \underline{d} \otimes \underline{e}_3 = \underline{e}_\alpha \otimes \underline{r}_{,\alpha} \underline{d} \otimes \underline{e}_3$$

$$= \underline{d} \cdot \underline{r}_{,\alpha} \underline{e}_\alpha \otimes \underline{e}_3 \equiv \underline{\gamma}_\alpha \underline{e}_\alpha \otimes \underline{e}_3$$

$$\underline{e}_3 \otimes \underline{d} \underline{\nabla}_r = \underline{e}_3 \otimes \underline{d} \underline{r}_{,\alpha} \otimes \underline{e}_\alpha = \underline{\gamma}_\alpha \underline{e}_3 \otimes \underline{e}_\alpha$$

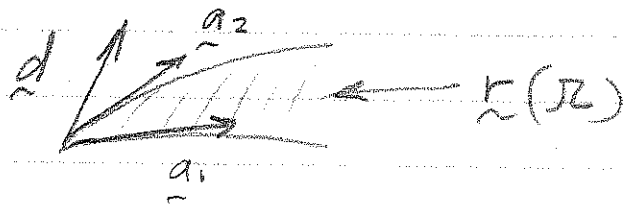
$$\begin{aligned} \underline{e}_3 \otimes \underline{d} (\underline{d} \otimes \underline{e}_3) &= |\underline{d}|^2 \underline{e}_3 \otimes \underline{e}_3 \\ &\equiv \lambda^2 \underline{e}_3 \otimes \underline{e}_3 \end{aligned}$$

Also,

$$\begin{aligned} \underline{F}_m^T \underline{G} &= (\underline{e}_3 \otimes \underline{d} + \nabla \underline{r}^T) (\underline{d}_{,\alpha} \otimes \underline{e}_\alpha) \\ &= \underline{d} \cdot \underline{d}_{,\alpha} \underline{e}_3 \otimes \underline{e}_\alpha + \underline{e}_\beta \otimes \underline{r}_{,\beta} \underline{d}_{,\alpha} \otimes \underline{e}_\alpha \\ &= \delta_\alpha \underline{e}_3 \otimes \underline{e}_\alpha + \underline{r}_{,\beta} \cdot \underline{d}_{,\alpha} \underline{e}_\beta \otimes \underline{e}_\alpha \\ &\quad + K_{\alpha\beta} \underline{e}_\alpha \otimes \underline{e}_\beta \end{aligned}$$

$$\delta_\alpha \equiv \underline{d} \cdot \underline{d}_{,\alpha}, \quad K_{\alpha\beta} \equiv \underline{r}_{,\alpha} \cdot \underline{d}_{,\beta}$$

Note that $\underline{a}_\alpha \equiv \underline{r}_{,\alpha}$ are tangent vectors on the surface $\underline{r}(\mathcal{J})$



Clearly $\delta_\alpha \equiv \underline{a}_\alpha \cdot \underline{d}$ are nonzero when \underline{d} is not \perp to the tangent plane $\text{span}\{\underline{a}_1, \underline{a}_2\}$, and we call δ_α the transverse shears.