

Let's examine (\*) p. 234 in detail:

$$\int_{\Omega} \frac{\partial \Xi}{\partial d_{i,\alpha}} \cdot \xi_{i,\alpha} dA = \int_{\Omega} \left( \underline{M}_1 \cdot \frac{\partial \xi}{\partial X_1} + \underline{M}_2 \cdot \frac{\partial \xi}{\partial X_2} \right) dA$$

Along edge  $E$ , it is the first term above that contributes to the boundary condition:

$$\int_0^L \int_0^L \underline{M}_1 \cdot \frac{\partial \xi}{\partial X_1} dX_1 dX_2$$

integrate by parts

$$\begin{aligned} \underline{U} &= \underline{M}_1 & d\underline{V} &= \frac{\partial \xi}{\partial X_1} dX_1 \\ d\underline{U} &= \frac{\partial \underline{M}_1}{\partial X_1} dX_1 & \underline{V} &= \xi \end{aligned}$$

$$= \int_0^L \left[ \underline{M}_1 \cdot \xi \Big|_{X_1=0}^{X_1=L} - \int_0^L \frac{\partial \underline{M}_1}{\partial X_1} \cdot \xi dX_1 \right] dX_2$$

Lecture 30

We want to examine the term at  $X_1=0$ :

$$(*) \int_0^L \underline{M}_1(0, X_2) \cdot \xi(0, X_2) dX_2 = 0 \quad \forall \text{ adm. variations } \xi$$

But we need to be careful here about admissibility. Returning to (\*) and (\*\*) p. 220, we need:

$$\begin{aligned} \frac{d}{d\varepsilon} \left[ (\underline{n} + \varepsilon \underline{\xi}) \cdot (\underline{n} + \varepsilon \underline{\xi}) = 1 \right] \Big|_{\varepsilon=0} \\ \Rightarrow \underline{n} \cdot \underline{\xi} = 0 \end{aligned}$$

Similarly,

$$\frac{d}{d\varepsilon} \left[ (\underline{a}_\alpha + \varepsilon \underline{n}_{1\alpha}) \cdot (\underline{n} + \varepsilon \underline{\xi}) = 0 \right] \Big|_{\varepsilon=0}$$

$$\Rightarrow \underline{a}_\alpha \cdot \underline{\xi} + \underline{n}_{1\alpha} \cdot \underline{n} = 0 \quad \alpha = 1, 2$$

$$\text{or } \boxed{\underline{a}_\alpha \cdot \underline{\xi} = -\underline{n}_{1\alpha} \cdot \underline{n}} \quad \alpha = 1, 2$$

Putting these two conditions together  $\Rightarrow$

$$(*) \quad \underline{\xi} = -(\underline{n}_{1\alpha} \cdot \underline{n}) \underline{a}^\alpha \quad \text{where } \{\underline{a}^1, \underline{a}^2\} \text{ is the}$$

reciprocal basis (so that  $\{\underline{a}^1, \underline{a}^2, \underline{n}\}$  is right-handed).

Thus, (\*) p. 235 becomes

$$\int_0^l \mu_1 \cdot \underline{a}^\alpha (\underline{n}_{1\alpha} \cdot \underline{n}) \Big|_{\underline{I}_1=0} d\underline{I}_2$$

Finally observe that admissibility for the geometric condition (2) p. 229 is

$$\underline{n}(0, \underline{I}_2) \equiv \underline{0} \quad \underline{I}_2 \in (0, l)$$

$$\Rightarrow \underline{n}_{12}(0, \underline{I}_2) \equiv \underline{0}$$

$$\rightarrow \therefore \int_0^l (\mu_1 \cdot \underline{a}^1) (\underline{n} \cdot \underline{n}_{11}) \Big|_{\underline{I}_1=0}^{d\underline{I}_2=0} \quad \forall \text{ adm. } \underline{n}_{11} \underline{\xi}$$

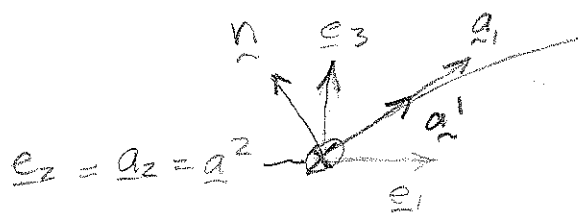
Since  $\underline{\xi}_1 = -\underline{n}_{11} \cdot \underline{n}$  is arbitrary, we conclude that

$$(*) \quad \underline{a}' \cdot \underline{M}_1 \Big|_{\underline{X}_1=0} = 0$$

(or we could prescribe this quantity as something non-zero.)

Observe from @ p. 234 that

$$\underline{a}_2(0, \underline{X}_2) = \underline{v}_{12}(0, \underline{X}_2) \equiv \underline{e}_2 = \underline{a}^2(0, \underline{X}_2)$$



Define  $\hat{\underline{a}}_1 = \frac{\underline{a}_1}{|\underline{a}_1|} = \frac{\underline{a}'}{|\underline{a}'|}$  at  $\underline{X}_1=0$ .

Then  $(*)$  above is equivalent to  $\hat{\underline{a}}_1 \cdot \underline{M}_1 \Big|_{\underline{X}_1=0} = 0$

Write  $\underline{M}_1 = (\hat{\underline{a}}_1 \cdot \underline{M}_1) \hat{\underline{a}}_1 + (\underline{e}_2 \cdot \underline{M}_1) \underline{e}_2 + (\underline{n} \cdot \underline{M}_1) \underline{n} \Big|_{\underline{X}_1=0}$

Then  $\underline{M}_1 \Big|_{\underline{X}_1=0} = \underline{n} \times \underline{M}_1 \Big|_{\underline{X}_1=0}$

$$= \left[ (\hat{\underline{a}}_1 \cdot \underline{M}_1) \underline{e}_2 - (\underline{e}_2 \cdot \underline{M}_1) \hat{\underline{a}}_1 \right] \Big|_{\underline{X}_1=0} = \left[ (\hat{\underline{a}}_1 \cdot \underline{M}_1) \hat{\underline{a}}_1 + (\underline{e}_2 \cdot \underline{M}_1) \underline{e}_2 \right] \Big|_{\underline{X}_1=0}$$

Hence  $(*)$  is equivalent to

$$(**) \quad \underline{e}_2 \cdot \underline{M}_1 \Big|_{\underline{X}_1=0} = 0$$



bending moment dist at  $x=1$

### ③ Free Edge:

In this case, no geometric b.c.'s are prescribed. Accordingly, there are two terms of interest coming from the 1st variation: (p. 222)

$$(*) \delta \tilde{\mathcal{F}} = \int_{\Omega} \left( \frac{\delta \tilde{\Phi}}{\delta \tilde{\tau}_{i,k}} + q_{\alpha} \tilde{n}_{\alpha} \right) \cdot \tilde{\tau}_{i,k} \\ + \frac{\delta \tilde{\Phi}}{\delta d_{i,\alpha}} \cdot \tilde{\tau}_{i,\alpha} + \dots$$

Next we integrate by parts (as before p. 235):

$$\int_0^L \int_0^L \left( \tilde{N}_1 + q_1 \tilde{n} \right) \cdot \frac{\delta \tilde{\tau}}{\delta \tilde{I}_1} + \dots + \tilde{\mu}_1 \cdot \frac{\delta \tilde{\tau}}{\delta \tilde{I}_1} + \dots d\tilde{I}_1 d\tilde{I}_2$$

$$= \int_0^L \left[ \left( \tilde{N}_1 + q_1 \tilde{n} \right) \cdot \tilde{\tau} + \tilde{\mu}_1 \cdot \tilde{\tau} \right] \Big|_{\tilde{I}_1=0}^{\tilde{I}_1=L} - \dots = 0$$

\(\forall \text{ adm } \tilde{\tau}, \tilde{\epsilon}\)

Once again, admissibility demands that we employ (\*) p. 236, which leads to

$$\int_0^L \left( \tilde{N}_1 + q_1 \tilde{n} \right) \cdot \tilde{\tau} - \left( \tilde{\mu}_1 \cdot \tilde{a}^{\alpha} \right) \left( \tilde{\tau}_{i,\alpha} \cdot \tilde{n} \right) \Big|_{\tilde{I}_1=0} d\tilde{I}_2 \\ = \int_0^L \left[ \left( \tilde{N}_1 + q_1 \tilde{n} \right) \cdot \tilde{\tau} - \left( \tilde{\mu}_1 \cdot \tilde{a}^1 \right) \tilde{n} \cdot \tilde{\tau}_{11} - \left( \tilde{\mu}_1 \cdot \tilde{a}^2 \right) \tilde{n} \cdot \frac{\delta \tilde{\tau}}{\delta \tilde{I}_2} \right] d\tilde{I}_2$$

As before  $\varepsilon_1 = -\underline{n} \cdot \underline{n}_{11}$  is arbitrary, and thus

$$(*) \quad \underline{a}^1 \cdot \underline{M}_1 \Big|_{\Sigma_1=0} = 0 \quad (\text{or prescribed})$$

We now integrate the third term in the integrand by parts - in this case the boundary terms are not of interest (influence other b.c.'s). We are then left with

$$\int_0^L \left[ (\underline{N}_1 + q_1 \underline{n}) \cdot \underline{n} + \frac{d}{dX_2} \left[ (\underline{a}^2 \cdot \underline{M}_1) \underline{n} \right] \cdot \underline{n} \right] d\text{ada } \underline{n} = 0$$

since  $\underline{n}$  is arbitrary, we conclude

$$(**) \quad \underline{N}_1 + q_1 \underline{n} + \frac{d}{dX_2} \left[ (\underline{a}^2 \cdot \underline{M}_1) \underline{n} \right] \Big|_{\Sigma_1=0} = 0 \quad \uparrow$$

or prescribed

Note that  $\frac{d}{dX_2} (\underline{a}^2 \cdot \underline{M}_1) \underline{n}$  combines with  $q_1 \underline{n}$  for the total "apparent" shear. This is the nonlinear version of Kirchhoff's famous result.

As before we can cast (\*) and (\*\*) in terms of the moment  $M_i$ :

Write

$$\underline{M}_1 = (\underline{a}^1 \cdot \underline{M}_1) \underline{a}_1 + (\underline{a}^2 \cdot \underline{M}_1) \underline{a}_2 + (\underline{n} \cdot \underline{M}_1) \underline{n}$$

Then  $\underline{M}_1 = \underline{n} \times \underline{M}_1$

$$= (\underline{a}^1 \cdot \underline{M}_1) \underline{n} \times \underline{a}_1 + (\underline{a}^2 \cdot \underline{M}_1) \underline{n} \times \underline{a}_2$$

It's not hard to see that

$$\underline{n} \times \underline{a}_1 = J \underline{a}^2 \quad J \equiv |\underline{a}_1 \times \underline{a}_2| \quad (\text{p. 220})$$

and

$$\underline{n} \times \underline{a}_2 = -J \underline{a}^1$$

(Indeed,  $\underline{a}^1 \cdot \underline{a}_2 = \underline{a}^2 \cdot \underline{a}_1 = 0$ , and

$$\underline{a}^1 \cdot \underline{a}_1 = \frac{1}{J} (\underline{a}_2 \times \underline{n}) \cdot \underline{a}_1 = \frac{1}{J} (\underline{a}_1 \times \underline{a}_2) \cdot \underline{n} = 1 \quad \checkmark$$


$$\underline{a}^2 \cdot \underline{a}_2 = \frac{1}{J} (\underline{n} \times \underline{a}_1) \cdot \underline{a}_2 = \frac{1}{J} (\underline{a}_1 \times \underline{a}_2) \cdot \underline{n} = 1 \quad \checkmark$$

$$\therefore \underline{M}_1 = (\underline{a}^1 \cdot \underline{M}_1) J \underline{a}^2 - (\underline{a}^2 \cdot \underline{M}_1) J \underline{a}^1$$

$$\Rightarrow \begin{cases} \underline{a}_1 \cdot \underline{M}_1 = -J \underline{a}^2 \cdot \underline{M}_1 \\ \underline{a}_2 \cdot \underline{M}_1 = J \underline{a}^1 \cdot \underline{M}_1 \end{cases}$$

So (\*)  $\Leftrightarrow$

$$\underline{a}_2 \cdot \underline{M}_1 \Big|_{x_1=0} = 0$$

or prescribed  
  
 prescribed bending moment at edge.

while (\*\*\*) is equivalent to

or prescribed

$$\left\{ \bar{N}_1 + q_1 \bar{n} - \frac{d}{dX_2} \left[ \left( \frac{a_1 \cdot M_1}{\bar{\sigma}} \right) \bar{n} \right] \right\} \Big|_{X_1=0} = 0$$

Note:  $q_1 - \frac{d}{dX_2} \left( \frac{a_1 \cdot M_1}{\bar{\sigma}} \right) \Big|_{X_1=0}$  is the apparent net shear distr.

while  $\left( \bar{N}_1 - \left( \frac{a_1 \cdot M_1}{\bar{\sigma}} \right) \frac{d\bar{n}}{dX_2} \right) \Big|_{X_1=0}$

is the apparent in-plane force distr.