

Lecture 13

This property (along with smoothness assumptions: smoothness of the boundary $\partial B = \partial B_1 \cup \partial B_2$ with $\partial B_1 \cap \partial B_2 = \emptyset$), yields a "local" existence result:

First re-write the equilibrium equations in terms of the displacement $\underline{x} + \underline{u}(\underline{x}) = \underline{f}(\underline{x})$
 $\Rightarrow \underline{F} = \underline{\underline{F}} + \nabla \underline{u}$, $\nabla^2 \underline{f} = \nabla^2 \underline{u}$:

$$(*) \quad F(\underline{\mu}, \underline{u}) = \begin{cases} \underline{\underline{C}}(\underline{\underline{F}} + \nabla \underline{u})[\nabla^2 \underline{u}] + \hat{\underline{b}} = \underline{0} & \text{in } B \\ \underline{u} - \hat{\underline{d}} = \underline{0} & \text{on } \partial B_1, \quad (\hat{\underline{d}} = \hat{\underline{f}} - \underline{x}) \\ \frac{dW}{d\underline{F}}(\underline{\underline{F}} + \nabla \underline{u}) \underline{m} - \hat{\underline{s}} = \underline{0} & \text{on } \partial B_2 \end{cases}$$

$\underline{\mu} \equiv (\hat{\underline{b}}, \hat{\underline{d}}, \hat{\underline{s}})$ prescribed fields.

Solve $F(\underline{\mu}, \underline{u}) = 0$ $F: \mathcal{P} \times \mathcal{X} \rightarrow \mathcal{Y}$
 $\mathcal{P}, \mathcal{X}, \mathcal{Y}$ (approp. Banach spaces)

Observe: $F(\underline{0}, \underline{0}) = \underline{0}$

$$\underline{\underline{C}}(\underline{\underline{F}}) \underline{0} + \hat{\underline{0}} = \underline{0} \quad \checkmark$$

$$\underline{0} - \underline{0} = \underline{0} \quad \checkmark$$

$$\frac{dW}{d\underline{F}}(\underline{\underline{F}}) \underline{0} - \underline{0} = \underline{0} \quad \checkmark$$

Recall implicit function theorem for finite-dimensional maps:

$$f(\mu, x) = 0 \quad f: \mathbb{R}^m \times \mathbb{R}^n \rightarrow \mathbb{R}^n$$

Assume: i) f is C^1

ii) $f(0, 0) = 0$ ($(0, 0)$ is a known solution)

iii) $\frac{\partial f}{\partial x}(0, 0) \in L(\mathbb{R}^n)$ is invertible.

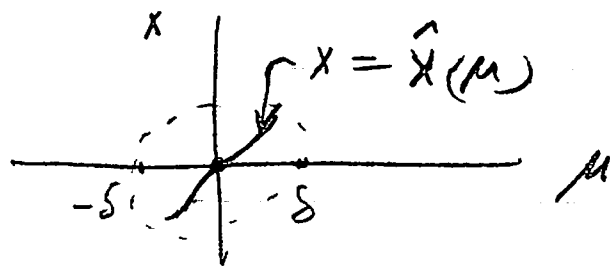
Then $x = \hat{x}(\mu)$ such that

(a) $f(\mu, \hat{x}(\mu)) = 0 \quad \forall |\mu| < \delta$

(b) $\hat{x}(0) = 0$

(c) \hat{x} is C^1

(d) "(a)" captures all local solutions



We can carry out the same procedure (formally) for (*) p. 103 :

We must linearize (in \underline{u}) about the $(\underline{u}, \underline{u}) = (0, \underline{0})$ solution:

$$\left. \frac{d}{d\alpha} F(0, \underline{0} + \alpha \underline{h}) \right|_{\alpha=0} \equiv D_u F(0, \underline{0}) [\underline{h}] \quad \forall \underline{h} \in \mathcal{X}$$

$$= \left. \frac{d}{d\alpha} \left[\underline{C}(\underline{\Gamma} + \alpha \underline{\nabla} \underline{h}) (\alpha \underline{\nabla}^2 \underline{h}) + \underline{0} \right] \right|_{\alpha=0}$$

$$\left. \frac{d}{d\alpha} \left[(\alpha \underline{h}) - \underline{0} \right] \right|_{\alpha=0}$$

$$\left. \frac{d}{d\alpha} \left[\frac{dW(\underline{\Gamma} + \alpha \underline{\nabla} \underline{h})}{d\underline{F}} \underline{m} - \underline{0} \right] \right|_{\alpha=0}$$

$$= \begin{cases} \underline{C}(\underline{\Gamma}) [\underline{\nabla}^2 \underline{h}] & \text{in } B \\ \underline{h} & \text{on } \partial B_1 \\ \underline{C}(\underline{\Gamma}) [\underline{\nabla} \underline{h}] \underline{m} & \text{on } \partial B_2 \end{cases}$$

Is the linear transformation $D_u F(0, \underline{0}) : \mathcal{X} \rightarrow \mathcal{Y}$ invertible?

1-1 : $D_u F(0, \underline{0}) [\underline{h}] = \underline{0}$ have a unique solution?

i.e., does

$$\begin{cases} \underline{C}(\underline{\Gamma}) \underline{\nabla}^2 \underline{h} = \underline{0} & \text{in } B \\ \underline{h} = \underline{0} & \text{on } \partial B_1 \\ \underline{C}(\underline{\Gamma}) [\underline{\nabla} \underline{h}] \underline{m} = \underline{0} & \text{on } \partial B_2 \end{cases}$$

have the unique solution $\underline{h} \equiv \underline{0}$?

If so, $D_h F(\underline{0}, \underline{0})$ is 1-1.

Ans: Yes

Pf: Kirchhoff's uniqueness theorem for linear elasticity. Observe

$$\underline{C}(\underline{\epsilon}) \underline{\nabla}^2 \underline{h} = \underline{\nabla} \cdot (\underline{C}(\underline{\epsilon}) [\underline{\nabla} \underline{h}])$$

"Mult." by \underline{h} and integrate over \mathcal{B} :

$$\int_{\mathcal{B}} \underline{h} \cdot \underline{\nabla} \cdot (\underline{C}(\underline{\epsilon}) [\underline{\nabla} \underline{h}]) dV$$

Again, from p. 35: $(\underline{\nabla} \cdot \underline{\xi}) \cdot \underline{n} = \underline{\nabla} \cdot (\underline{\xi}^T \underline{n})$

$$\underline{\xi} \equiv \underline{C}(\underline{\epsilon}) [\underline{\nabla} \underline{h}] \quad - \underline{\xi} \cdot \underline{\nabla} \underline{n}$$

div
thm
=

$$- \int_{\mathcal{B}} \underline{\nabla} \underline{h} \cdot \underline{C}(\underline{\epsilon}) [\underline{\nabla} \underline{h}] dV$$

$$+ \int_{\partial \mathcal{B}_2} \underline{C}(\underline{\epsilon}) [\underline{\nabla} \underline{h}] \cdot \underline{n} \cdot \underline{h} \rightarrow 0 \text{ b.c.}$$

$$\therefore \int_{\mathcal{B}} \underline{\nabla} \underline{h} \cdot \underline{C}(\underline{\epsilon}) [\underline{\nabla} \underline{h}] dV = 0$$

Recall $\underline{C}(\underline{\epsilon}) [\underline{W}] = \underline{0} \quad \forall \underline{W} \in \text{skew}$.

For a given $\underline{\nabla}u$, write $\underline{\nabla}u = \underline{H} + \underline{W}$,
 where $\underline{H} = \frac{1}{2}(\underline{\nabla}u + \underline{\nabla}u^T)$ and $\underline{W} = \frac{1}{2}(\underline{\nabla}u - \underline{\nabla}u^T)$.

$$\begin{aligned} \text{Then } & (\underline{H} + \underline{W}) \cdot \underline{C}(\underline{\mathbb{E}}) [\underline{H} + \underline{W}] \\ &= (\underline{H} + \underline{W}) \cdot \underline{C}(\underline{\mathbb{E}}) [\underline{H}] \quad (\underline{C}(\underline{\mathbb{E}}) [\underline{W}] = \underline{0}) \\ &= \underline{H} \cdot \underline{C}(\underline{\mathbb{E}}) [\underline{H}] + \underline{C}(\underline{\mathbb{E}}) [\underline{W}] \cdot \underline{H} \quad (\text{Lemma p. 99}) \\ &> 0 \quad \forall \underline{H} \in \text{Sym}(\mathbb{E}^3) \\ &\quad \underline{H} \neq \underline{0}. \end{aligned}$$

Thus we conclude that $\underline{\nabla}u(\underline{x}) \in \text{Skew}(\mathbb{E}^3)$.
 Now if u is C^2 (which we assume), then

$$\underline{\nabla}u = -\underline{\nabla}u^T$$

$$\Rightarrow \frac{\partial u_i}{\partial x_j} = -\frac{\partial u_j}{\partial x_i}$$

$$\text{Then } \frac{\partial^2 u_i}{\partial x_k \partial x_j} = -\frac{\partial^2 u_j}{\partial x_k \partial x_i}$$

$$\Rightarrow \frac{\partial}{\partial x_j} \left(\frac{\partial u_i}{\partial x_k} \right) = -\frac{\partial}{\partial x_i} \left(\frac{\partial u_j}{\partial x_k} \right)$$

$$\Rightarrow -\frac{\partial}{\partial x_j} \left(\frac{\partial u_k}{\partial x_i} \right) = \frac{\partial}{\partial x_i} \left(\frac{\partial u_k}{\partial x_j} \right) \Rightarrow \frac{\partial^2 u_k}{\partial x_i \partial x_j} = 0$$

$$\Rightarrow \underline{\nabla}(\underline{\nabla}u) = \underline{0}$$

$$\Rightarrow \underline{\nabla}u(\underline{x}) = \underline{W}_0 \in \text{Skew}$$

$$\Rightarrow \underline{u} = \underline{W}_0 \underline{x} + \underline{c}$$

$$\underline{u}|_{\partial B_1} = \underline{0} \Rightarrow \underline{W}_0 = \underline{0}, \underline{c} = \underline{0}.$$

To show that $D_u F(0, \underline{e})$ is onto, we would need to demonstrate that any inhomogeneous problem $D_u F(0, \underline{e})[\underline{h}] = \hat{\underline{g}}$ has a solution, viz.)

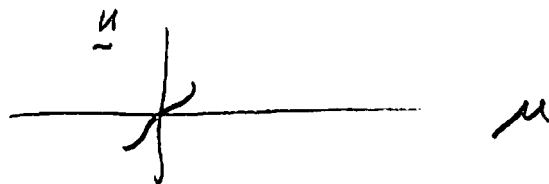
$$\left\{ \begin{array}{l} \mathcal{L}(\underline{x}) \nabla^2 \underline{h} = \hat{\underline{g}}_1 \quad \text{in } B, \\ \underline{h} = \hat{\underline{g}}_2 \quad \text{on } \partial B_1, \\ \mathcal{L}(\underline{x}) [\nabla \underline{h}]_{\underline{m}} = \hat{\underline{g}}_3 \quad \text{on } \partial B_2, \end{array} \right.$$

has a solution " \underline{h} ". This is much harder - but doable - using elliptic pde theory (which we must skip here).

With this in hand, we then have a "path" of solutions

$$\underline{u}(\underline{x}) = \hat{\underline{u}}(\underline{x}; (\hat{\underline{b}}, \hat{\underline{d}}, \hat{\underline{a}})^{\mu})$$

for " $|\mu|$ sufficiently small", to the nonlinear problem.



To motivate our next hypothesis for $\underline{C}(\underline{F})$, we suppose that $\underline{f}_0 = \underline{X} + \underline{u}_0$ is an equilibrium solution, viz.,

$$\underline{C}(\underline{F} + \underline{u}_0) [\nabla_{\underline{u}_0}^2] + \underline{\hat{b}} \equiv \underline{0} \text{ in } \mathcal{B},$$

$$\underline{u}_0 \equiv \underline{\hat{d}} \text{ on } \downarrow \mathcal{B}_1,$$

$$\frac{dW}{d\underline{F}}(\underline{F}_0) \underline{m} \equiv \underline{\hat{J}} \text{ on } \downarrow \mathcal{B}_2,$$

and we linearize the dynamical equation: (p. 77) at \underline{u}_0 :

$$\underline{f}(\underline{X}, t) = \underline{f}_0(\underline{X}) + \alpha \underline{w}(\underline{X}, t)$$

Plug into dynamical equations and linearize:

$$\frac{d}{d\alpha} \left[\underline{C}(\underline{F}_0 + \alpha \underline{\nabla} \underline{w}) [\underline{\nabla}^2 \underline{f}_0 + \alpha \underline{\nabla}^2 \underline{w}] + \underline{\hat{b}} \right. \\ \left. = \underline{f}_0(\alpha \underline{w}_{tt}) \right] \Big|_{\alpha=0}$$

$$\frac{d}{d\alpha} \left(\underline{f}_0 + \alpha \underline{w} - \underline{\hat{f}} \right) \Big|_{\alpha=0} = \underline{0}$$

$$\frac{d}{d\alpha} \left(\frac{dW}{d\underline{F}}(\underline{F}_0 + \alpha \underline{\nabla} \underline{w}) \underline{m} - \underline{\hat{J}} \right) \Big|_{\alpha=0} = \underline{0},$$

which leads to:

$$\left\{ \begin{array}{l}
 \underline{C}(\underline{F}_0) [\underline{\nabla}^2 \underline{w}] + \overbrace{\frac{d\underline{C}(\underline{F}_0)}{d\underline{F}}(\underline{F}_0) [\underline{\nabla} \underline{w}]}^{\text{lower order}} \underline{\nabla}^2 \underline{f}_0 \\
 \\
 = \rho_0 \underline{w}_{tt} \quad \text{in } B \\
 \\
 \underline{w} = \underline{0} \quad \text{on } \partial B_1 \\
 \\
 \underline{C}(\underline{F}_0) [\underline{\nabla} \underline{w}] \underline{n} = \underline{0} \quad \text{on } \partial B_2
 \end{array} \right.$$

Linearized dynamical equations about an arbitrary equilibrium. Let's make the assumption that \underline{f}_0 is homogeneous, i.e., $\underline{F}_0 = \underline{\text{const}}$. Then the pde becomes:

$$(*) \quad \underline{\nabla} \cdot (\underline{C}(\underline{F}_0) [\underline{\nabla} \underline{w}]) = \rho_0 \underline{w}_{tt}$$

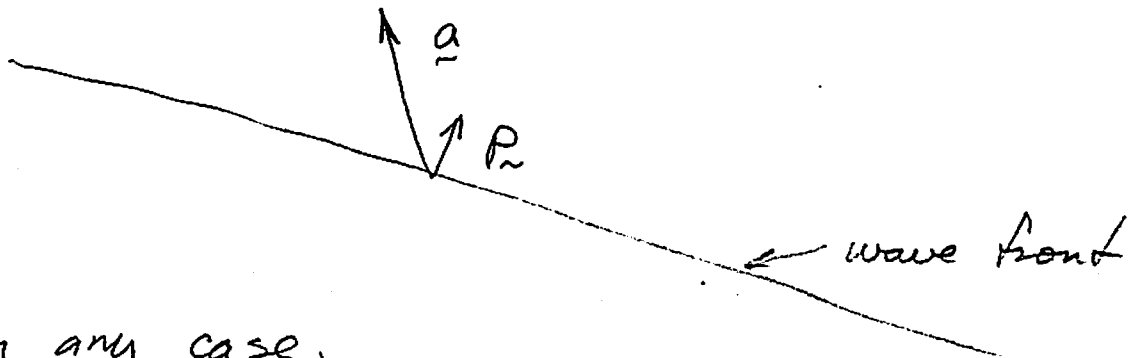
Let's also ignore the b.c.'s (eg. infinite medium) and use wave propagation to guide us in a proper assumption for $\underline{C}(\underline{F}_0)$ (recall, pos-def'n on sym. tensors is not good for $\underline{F}_0 \neq \underline{I}$, pp. 94-96)

Seek a planar travelling wave solution of (*):

$$\underline{w}(\underline{x}, t) = \underline{a} \, g(\underline{p} \cdot \underline{x} - ct),$$

where \underline{p} is a unit vector \perp to the

wave front; \underline{p} is called the direction of propagation, and \underline{a} is the direction of motion:



In any case,

$$\begin{aligned} \underline{\nabla} \cdot \underline{C}(\underline{F}_0) \underline{\nabla} w &= \frac{d}{dt} \underline{a} \cdot \underline{g}(\underline{p} \cdot (\underline{x} + \alpha \underline{n}) - ct) \Big|_{d=0} \\ &= \underline{a} \cdot \underline{g}'(\underline{p} \cdot \underline{x} - ct) \underline{p} \cdot \underline{n} \\ &= \underbrace{[\underline{a} \otimes \underline{p} \cdot \underline{g}'(\underline{p} \cdot \underline{x} - ct)]}_{\underline{\nabla} w} \cdot \underline{n} \quad \forall \underline{n} \in \mathbb{E}^3 \end{aligned}$$

$$\text{Then } \left(\underline{\nabla} \cdot \left(\underline{C}(\underline{F}_0) [\underline{\nabla} w] \right) \right)_i$$

$$= \left(\underline{\nabla} \cdot \left(\underline{C}(\underline{F}_0) \underbrace{[\underline{a} \otimes \underline{p}]}_{\underline{B}^0} \cdot \underline{g}'(\underline{p} \cdot \underline{x} - ct) \right) \right)_i$$

$$= \frac{1}{\underline{J} \underline{x}_j} \left(B_{ij}^0 \cdot \underline{g}'(p_k \underline{x}_k - ct) \right)$$

$$= B_{ij}^0 \cdot \underline{g}''(p_k \underline{x}_k - ct) p_k \delta_{kj}^T p_j$$

$$\nabla \cdot (\underline{C}(\underline{F}_0) [\underline{v} \underline{w}])$$

$$= g''(\cdot) \underline{C}(\underline{F}_0) [\underline{a} \otimes \underline{p}] \underline{p}$$

$\therefore (*) \Rightarrow$

$$g''(\cdot) \underline{C}(\underline{F}_0) [\underline{a} \otimes \underline{p}] \underline{p} = \rho_0 c^2 \underline{a} g''(\cdot)$$

(assuming $g''(\cdot) \neq 0$)

$$\Rightarrow \boxed{\underline{C}(\underline{F}_0) [\underline{a} \otimes \underline{p}] \underline{p} = \rho_0 c^2 \underline{a}} \quad (*)$$

question: for fixed \underline{p} , why does this define a tensor?

Defn $\underline{Q}(\underline{p}) \underline{a} \equiv \underline{C}(\underline{F}_0) [\underline{a} \otimes \underline{p}] \underline{p}$, where

$\underline{Q}(\underline{p})$ is called the acoustic tensor in the direction of \underline{p} .

For a given \underline{p} , $(*)$ above is an eigenvalue problem:

$$\boxed{\underline{Q}(\underline{p}) \underline{a} = \rho_0 c^2 \underline{a}} \quad (**)$$

Observe:

$$\begin{aligned} \underline{b} \cdot \underline{Q}(\underline{p}) \underline{a} &= \underline{b} \cdot \underline{C}(\underline{F}_0) [\underline{a} \otimes \underline{p}] \underline{p} \\ &= b_i C_{ijkl} a_k p_l p_j \\ &= b_i p_j C_{ijkl} a_k p_l \\ &= \underline{b} \otimes \underline{p} \cdot \underline{C}(\underline{F}_0) [\underline{a} \otimes \underline{p}] \end{aligned}$$

Lemma p. 99

$$= \underline{a} \otimes \underline{r} \cdot \underline{C}(\underline{F}_0) [\underline{r} \otimes \underline{r}]$$

$$= \underline{a} \cdot \underline{Q}(\underline{r}) \underline{a},$$

i.e., $\underline{Q}^T = \underline{Q}$, so the eigenvalues

$\sigma = \rho_0 c^2$ of \underline{Q} are real. We would also like them to be positive (otherwise we could have imaginary wave speeds)

Require $\underline{a} \cdot \underline{Q}(\underline{r}) \underline{a} > 0$

or (SE) $\underline{a} \otimes \underline{r} \cdot \underline{C}(\underline{F}_0) [\underline{r} \otimes \underline{r}] > 0 \quad \forall \underline{a}, \underline{r} \neq \underline{0}$.

This is called strong ellipticity at \underline{F}_0 .
Insures well-posed "incremental" dynamics.

Exercise (18) For an isotropic solid we have

$$\underline{C}(\underline{F}) [\underline{H}] = \Lambda (\text{tr } \underline{H}) \underline{F} + 2\mu \underline{H} \quad \forall \underline{H} \in \text{Sym}(\mathbb{E}^3)$$

where Λ, μ are the Lamé constants (see any book on linear elasticity). (a) Show that the acoustic tensor (at the reference configuration, has the form

$$\underline{Q}(\underline{r}) = (\Lambda + 2\mu) \underline{r} \otimes \underline{r} + \mu (\underline{F} - \underline{r} \otimes \underline{r})$$

(b) Deduce: strong ellipticity of $\underline{C}(\underline{F}) \Leftrightarrow$

$$\Lambda + 2\mu > 0 \quad \text{and} \quad \mu > 0$$

(c) For a given direction of propagation " \underline{p} ", compute the possible wave speeds " c " (from the eigenvalues) and corresponding directions of motion " \underline{a} " (eigenvectors)

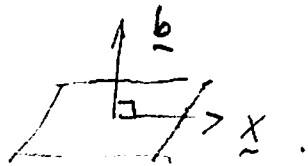
Note that (SE) is weaker than positive definiteness, which reads

$$(*) \quad \underline{A} \cdot \underline{C}(\underline{F}_0)[\underline{A}] > 0 \quad \forall \underline{A} \in L(\mathbb{E}^3), \underline{A} \neq \underline{0}.$$

Indeed, whereas (*) above is valid for any tensor $\underline{A} \neq \underline{0}$, (SE) involves only tensors of the form $\underline{a} \otimes \underline{b}$, $\underline{a}, \underline{b} \neq \underline{0}$, which always have rank = 1:

$$(\underline{a} \otimes \underline{b}) \underline{x} \equiv \underline{a} (\underline{b} \cdot \underline{x}) = \underline{0} \quad \forall \underline{x} \perp \underline{b},$$

ie., $\underline{a} \otimes \underline{b}$ has a 2-dimensional null space (the plane \perp to \underline{b}):

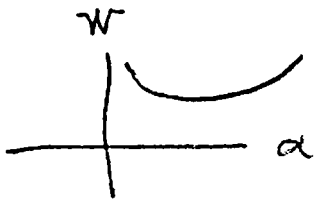


In fact, (*) \Rightarrow convexity, which as we know is not good physics. To see this, consider the function

$$\alpha \mapsto W(\underline{F}_0 + \alpha \underline{A}) \quad (\underline{F}_0, \underline{A} \neq \underline{0} \text{ fixed})$$

$$\text{Then } \frac{d^2}{d\alpha^2} W(\underline{F}_0 + \alpha \underline{A}) = \frac{d}{d\alpha} \left(\frac{dW}{d\underline{F}}(\underline{F}_0 + \alpha \underline{A}) \cdot \underline{A} \right)$$

positive 2nd derivative!



$$= \underline{A} \cdot \frac{d^2 W}{d\underline{F}^2}(\underline{F}_0 + \alpha \underline{A}) [\underline{A}] > 0$$

from (*),

$\Rightarrow \alpha \mapsto W(\underline{F}_0 + \alpha \underline{A})$ is convex - which is (can be shown) equivalent to convexity of $W(\underline{F})$ (p. 93).

On the other hand, (SE) also implies a kind of convexity, but only in "rank-one" directions, i.e., only "rank-one excursions" from a given \underline{F}_0 are allowed: $\underline{F}_0 + \alpha \underline{a} \otimes \underline{b}$

$\underline{a}, \underline{b} \neq \underline{0}$

$$\text{Then } \left. \frac{d}{d\alpha} W(\underline{F}_0 + \alpha \underline{a} \otimes \underline{b}) \right|_{\alpha=0} = \underline{a} \otimes \underline{b} \cdot \frac{dW}{d\underline{F}}(\underline{F}_0)$$

$$= \tilde{S}_{\alpha\beta}(\underline{F}_0) \text{ comp. 1st PK}$$

$$\text{Also } \left. \frac{d^2}{d\alpha^2} W(\underline{F}_0 + \alpha \underline{a} \otimes \underline{b}) \right|_{\alpha=0}$$

$$= \left. \frac{d}{d\alpha} \left[\frac{dW}{d\underline{F}}(\underline{F}_0 + \alpha \underline{a} \otimes \underline{b}) \cdot \underline{a} \otimes \underline{b} \right] \right|_{\alpha=0}$$

$$\equiv \frac{d \tilde{S}_{\alpha\beta}(\underline{F}_0)}{dF_{\alpha\beta}} = \underline{a} \otimes \underline{b} \cdot \underline{C}(\underline{F}_0) [\underline{a} \otimes \underline{b}] > 0$$